

## **Supporting Information**

### **Quantum Mechanical Theory of Dynamic Nuclear Polarization in Solid Dielectrics**

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## Section I – Symbols used in the text

- $\delta$  - homogeneous line width of EPR spectrum
- $\Delta$  - inhomogeneous breadth of EPR spectrum
- $\omega_{1S}$  - strength of the microwave irradiation field
- $\omega_M$  - microwave irradiation frequency
- $\omega_{0S_i}$  - Larmor frequency of electron  $i$
- $\omega_0$  - nuclear Larmor frequency
- $\gamma_S$  - electron gyromagnetic ratio
- $\gamma_I$  - nuclear gyromagnetic ratio
- $B_0$  - static magnetic field

## Section II – Commutator relations (Eq.44)

1)  $M_{\Delta x}$  commutes with both  $S_{\Delta z}I_z$  and  $S_{1z}S_{2z}$  since

$$\begin{aligned} S_{\Delta z}I_z &= \frac{1}{4}(E_{M\Sigma} - E_{M\Delta}) \\ S_{1z}S_{2z} &= \frac{1}{4}(E_{S\Sigma} - E_{M\Sigma} - E_{M\Delta}) \end{aligned} \quad (\text{S1})$$

2) From the definitions in Table 1, it follows that

$$S_{\Delta}^{\pm}S_{\Sigma z} = S_1^{\pm}S_2^{\mp} \cdot \frac{1}{2}(S_{1z} + S_{2z}) = 0 \quad (\text{S2})$$

Therefore,

$$M_{\Delta x}S_{\Sigma z} = \frac{1}{2}(S_{\Delta}^+I^- + S_{\Delta}^-I^+)S_{\Sigma z} = 0 \quad (\text{S3})$$

3) Based on the results of Eq. S3, it also follows that

$$[M_{\Delta x}, S_{\Sigma z}I_z] = M_{\Delta x}S_{\Sigma z}I_z - S_{\Sigma z}I_zM_{\Delta x} = 0 \quad (\text{S4})$$

$$[M_{\Delta x}, E_{S\Sigma}I_z] = M_{\Delta x}E_{S\Sigma}I_z - E_{S\Sigma}I_zM_{\Delta x} = 4(M_{\Delta x}S_{\Sigma z}^2I_z - S_{\Sigma z}^2I_zM_{\Delta x}) = 0 \quad (\text{S5})$$

4) Going back to the definitions in Table 1, the following commutator relation can also be obtained:

$$\begin{aligned} [M_{\Delta x}, M_{\Sigma z}] &= \frac{1}{4}[S_{\Delta}^+I^- + S_{\Delta}^-I^+, S_{\Delta z} + E_{S\Delta}I_z] \\ &= \frac{1}{4}(-S_{\Delta}^+I^- + S_{\Delta}^-I^+ + S_{\Delta}^+I^- - S_{\Delta}^-I^+) \\ &= 0. \end{aligned} \quad (\text{S6})$$

### Section III – The CE Microwave Hamiltonian in the EBS

To obtain the effective microwave excitations for DNP, the microwave Hamiltonian  $\tilde{H}_M = 2\omega_{1S}(S_{1x} + S_{2x})\cos(\omega_M t)$  for the two electrons needs to be transformed from the PSB to the EBS using the unitary transformations  $U_\zeta$  (Eq. 33) and  $U_\xi$  (Eq.45):

$$\tilde{H}_M = 2\omega_{1S}\cos(\omega_M t) \cdot U_\xi U_\zeta (S_{1x} + S_{2x}) U_\zeta^{-1} U_\xi^{-1} \quad (\text{S7} = 51)$$

First, we evaluate the effect of  $U_\zeta$ , which is given by the following expression:

$$U_\zeta = \exp[i\zeta_\alpha S_{\Delta y} I^\alpha + i\zeta_\beta S_{\Delta y} I^\beta], \quad (\text{S8} = 33)$$

Let,  $U_\phi = e^{i\phi S_{\Delta y}} = \exp[\frac{1}{2}\phi(S_1^+ S_2^- - S_1^- S_2^+)]$ , where  $\phi$  represents  $\zeta_\alpha$  or  $\zeta_\beta$  in Eq. S8.

Then,

$$\begin{aligned} U_\phi S_{1z} U_\phi^{-1} &= U_\phi (S_{\Sigma z} + S_{\Delta z}) U_\phi^{-1} \\ &= S_{\Sigma z} + S_{\Delta z} \cos \phi - S_{\Delta x} \sin \phi \\ &= \frac{1}{2}(S_{1z} + S_{2z}) + \frac{1}{2}(S_{1z} - S_{2z}) \cos \phi - \frac{1}{2}(S_1^+ S_2^- + S_1^- S_2^+) \sin \phi, \end{aligned} \quad (\text{S9})$$

$$\begin{aligned} U_\phi S_1^+ U_\phi^{-1} &= S_1^+ + \frac{\phi}{2}(2S_{1z} S_2^+) + \left(\frac{\phi}{2}\right)^2 \frac{1}{2!}(-S_1^+) + \left(\frac{\phi}{2}\right)^3 \frac{1}{3!}(-2S_{1z} S_2^+) + \dots \\ &= S_1^+ \cos \frac{\phi}{2} + 2S_{1z} S_2^+ \sin \frac{\phi}{2}, \end{aligned} \quad (\text{S10})$$

$$\begin{aligned} U_\phi S_1^- U_\phi^{-1} &= S_1^- + \frac{\phi}{2}(2S_{1z} S_2^-) + \left(\frac{\phi}{2}\right)^2 \frac{1}{2!}(-S_1^-) + \left(\frac{\phi}{2}\right)^3 \frac{1}{3!}(-2S_{1z} S_2^-) + \dots \\ &= S_1^- \cos \frac{\phi}{2} + 2S_{1z} S_2^- \sin \frac{\phi}{2}. \end{aligned} \quad (\text{S11})$$

Similarly,

$$\begin{aligned} U_\phi S_{2z} U_\phi^{-1} &= U_\phi (S_{\Sigma z} - S_{\Delta z}) U_\phi^{-1} \\ &= S_{\Sigma z} - S_{\Delta z} \cos \phi + S_{\Delta x} \sin \phi \\ &= \frac{1}{2}(S_{1z} + S_{2z}) - \frac{1}{2}(S_{1z} - S_{2z}) \cos \phi + \frac{1}{2}(S_1^+ S_2^- + S_1^- S_2^+) \sin \phi, \end{aligned} \quad (\text{S12})$$

$$\begin{aligned} U_\phi S_2^+ U_\phi^{-1} &= S_2^+ + \frac{\phi}{2}(-2S_1^+ S_{2z}) + \left(\frac{\phi}{2}\right)^2 \frac{1}{2!}(-S_2^+) + \left(\frac{\phi}{2}\right)^3 \frac{1}{3!}(2S_1^+ S_{2z}) + \dots \\ &= S_2^+ \cos \frac{\phi}{2} - 2S_1^+ S_{2z} \sin \frac{\phi}{2}, \end{aligned} \quad (\text{S13})$$

$$\begin{aligned} U_\phi S_2^- U_\phi^{-1} &= S_2^- + \frac{\phi}{2} (-2S_1^- S_{2z}) + \left(\frac{\phi}{2}\right)^2 \frac{1}{2!} (-S_2^-) + \left(\frac{\phi}{2}\right)^3 \frac{1}{3!} (2S_1^- S_{2z}) + \dots \\ &= S_2^- \cos \frac{\phi}{2} - 2S_1^- S_{2z} \sin \frac{\phi}{2}, \end{aligned} \quad (\text{S14})$$

Combining these results, leads to Eq. 52:

$$\begin{aligned} \tilde{H}_M &= U_\zeta (S_{1x} + S_{2x}) U_\zeta^{-1} \\ &= S_{1x} (I^\alpha \cos \frac{\zeta_\alpha}{2} + I^\beta \cos \frac{\zeta_\beta}{2}) + 2S_{1z} S_{2x} (I^\alpha \sin \frac{\zeta_\alpha}{2} + I^\beta \sin \frac{\zeta_\beta}{2}) \\ &\quad + S_{2x} (I^\alpha \cos \frac{\zeta_\alpha}{2} + I^\beta \cos \frac{\zeta_\beta}{2}) - 2S_{1x} S_{2z} (I^\alpha \sin \frac{\zeta_\alpha}{2} + I^\beta \sin \frac{\zeta_\beta}{2}) \\ &= \frac{1}{2} (S_{1x} + S_{2x}) (c_\alpha + c_\beta) + (S_{1x} + S_{2x}) I_z (c_\alpha - c_\beta) \\ &\quad + (S_{1z} S_{2x} - S_{1x} S_{2z}) (s_\alpha + s_\beta) + 2(S_{1z} S_{2x} - S_{1x} S_{2z}) I_z (s_\alpha - s_\beta), \end{aligned} \quad (\text{S15} = 52)$$

$$c_\alpha = \cos \frac{\zeta_\alpha}{2}, c_\beta = \cos \frac{\zeta_\beta}{2}, s_\alpha = \sin \frac{\zeta_\alpha}{2}, s_\beta = \sin \frac{\zeta_\beta}{2}. \quad (\text{S16} = 53)$$

Next, we evaluate the effect of  $U_\xi$ , which is given in Eq. 45, i.e.

$$U_\xi = \exp[i\xi M_{\Delta y}], \quad (\text{S17} = 45)$$

Let  $U_\chi = e^{i\chi M_{\Delta y}} = \exp[\frac{1}{2}\chi(S_1^+ S_2^- I^- - S_1^- S_2^+ I^+)]$ , where  $\chi$  represents  $\xi$  in Eq. 45. Then we can derive the following intermediate expressions:

$$\begin{aligned} U_\chi I_z U_\chi^{-1} &= U_\chi (E_{S\Sigma} I_z + E_{S\Delta} I_z) U_\chi^{-1} = U_\chi (E_{S\Sigma} I_z + M_{\Sigma z} - M_{\Delta z}) U_\chi^{-1} \\ &= E_{S\Sigma} I_z + M_{\Sigma z} - M_{\Delta z} \cos \chi + M_{\Delta z} \sin \chi \\ &= \frac{1}{4} S_{1z} - \frac{1}{4} S_{2z} + \frac{3}{4} I_z + S_{1z} S_{2z} I_z \\ &\quad - \left(\frac{1}{4} S_{1z} - \frac{1}{4} S_{2z} - \frac{1}{4} S_{2z} + S_{1z} S_{2z} I_z\right) \cos \chi \\ &\quad + \frac{1}{2} (S_1^+ S_2^- I^- + S_1^- S_2^+ I^+) \sin \chi, \end{aligned} \quad (\text{S18})$$

$$\begin{aligned} U_\chi I^+ U_\chi^{-1} &= I^+ + \frac{\chi}{2} (-2S_1^+ S_2^- I_z) + \left(\frac{\chi}{2}\right)^2 \frac{1}{2!} \left(\frac{-1}{2} I^+ + 2S_{1z} S_{2z} I^+\right) \\ &\quad + \left(\frac{\chi}{2}\right)^3 \frac{1}{3!} (2S_1^+ S_2^- I_z) + \left(\frac{\chi}{2}\right)^4 \frac{1}{4!} \left(\frac{1}{2} I^+ - 2S_{1z} S_{2z} I^+\right) + \dots \\ &= \frac{1}{2} I^+ + 2S_{1z} S_{2z} I^+ \\ &\quad + \left(\frac{1}{2} I^+ - 2S_{1z} S_{2z} I^+\right) \cos \frac{\chi}{2} - 2S_1^+ S_2^- I_z \sin \frac{\chi}{2}, \end{aligned} \quad (\text{S19})$$

$$\begin{aligned}
U_\chi I^- U_\chi^{-1} &= I^- + \frac{\chi}{2} (-2S_1^- S_2^+ I_z) + \left(\frac{\chi}{2}\right)^2 \frac{1}{2!} \left(\frac{-1}{2} I^- + 2S_{1z} S_{2z} I^-\right) \\
&\quad + \left(\frac{\chi}{2}\right)^3 \frac{1}{3!} (2S_1^- S_2^+ I_z) + \left(\frac{\chi}{2}\right)^4 \frac{1}{4!} \left(\frac{1}{2} I^- - 2S_{1z} S_{2z} I^-\right) + \dots \\
&= \frac{1}{2} I^- + 2S_{1z} S_{2z} I^- \\
&\quad + \left(\frac{1}{2} I^- - 2S_{1z} S_{2z} I^-\right) \cos \frac{\chi}{2} - 2S_1^- S_2^+ I_z \sin \frac{\chi}{2}.
\end{aligned} \tag{S20}$$

Further,

$$\begin{aligned}
U_\chi S_{1z} U_\chi^{-1} &= U_\chi (S_{\Sigma z} + M_{\Sigma z} + M_{\Delta z}) U_\chi^{-1} \\
&= S_{\Sigma z} + M_{\Sigma z} + M_{\Delta z} \cos \chi - M_{\Delta x} \sin \chi \\
&= \frac{3}{4} S_{1z} + \frac{1}{4} S_{2z} + \frac{1}{4} I_z - S_{1z} S_{2z} I_z \\
&\quad + \left(\frac{1}{4} S_{1z} - \frac{1}{4} S_{2z} - \frac{1}{4} S_{2z} + S_{1z} S_{2z} I_z\right) \cos \chi \\
&\quad - \frac{1}{2} (S_1^+ S_2^- I^- + S_1^- S_2^+ I^+) \sin \chi,
\end{aligned} \tag{S21}$$

$$\begin{aligned}
U_\chi S_1^+ U_\chi^{-1} &= S_1^+ + \frac{\chi}{2} (2S_{1z} S_2^+ I^+) + \left(\frac{\chi}{2}\right)^2 \frac{1}{2!} \left(\frac{-1}{2} S_1^+ - 2S_1^+ S_{2z} I_z\right) \\
&\quad + \left(\frac{\chi}{2}\right)^3 \frac{1}{3!} (-2S_{1z} S_2^+ I^+) + \left(\frac{\chi}{2}\right)^4 \frac{1}{4!} \left(\frac{1}{2} S_1^+ + 2S_1^+ S_{2z} I_z\right) + \dots \\
&= \frac{1}{2} S_1^+ - 2S_1^+ S_{2z} I_z \\
&\quad + \left(\frac{1}{2} S_1^+ + 2S_1^+ S_{2z} I_z\right) \cos \frac{\chi}{2} + 2S_{1z} S_2^+ I^+ \sin \frac{\chi}{2},
\end{aligned} \tag{S22}$$

$$\begin{aligned}
U_\chi S_1^- U_\chi^{-1} &= S_1^- + \frac{\chi}{2} (2S_{1z} S_2^- I^-) + \left(\frac{\chi}{2}\right)^2 \frac{1}{2!} \left(\frac{-1}{2} S_1^- - 2S_1^- S_{2z} I_z\right) \\
&\quad + \left(\frac{\chi}{2}\right)^3 \frac{1}{3!} (-2S_{1z} S_2^- I^-) + \left(\frac{\chi}{2}\right)^4 \frac{1}{4!} \left(\frac{1}{2} S_1^- + 2S_1^- S_{2z} I_z\right) + \dots \\
&= \frac{1}{2} S_1^- - 2S_1^- S_{2z} I_z \\
&\quad + \left(\frac{1}{2} S_1^- + 2S_1^- S_{2z} I_z\right) \cos \frac{\chi}{2} + 2S_{1z} S_2^- I^- \sin \frac{\chi}{2}.
\end{aligned} \tag{S23}$$

Finally,

$$\begin{aligned}
U_\chi S_{2z} U_\chi^{-1} &= U_\chi (S_{\Sigma z} - M_{\Sigma z} - M_{\Delta z}) U_\chi^{-1} \\
&= S_{\Sigma z} - M_{\Sigma z} - M_{\Delta z} \cos \chi + M_{\Delta x} \sin \chi \\
&= \frac{1}{4} S_{1z} + \frac{3}{4} S_{2z} - \frac{1}{4} I_z + S_{1z} S_{2z} I_z \\
&\quad - \left(\frac{1}{4} S_{1z} - \frac{1}{4} S_{2z} - \frac{1}{4} S_{2z} + S_{1z} S_{2z} I_z\right) \cos \chi \\
&\quad + \frac{1}{2} (S_1^+ S_2^- I^- + S_1^- S_2^+ I^+) \sin \chi.
\end{aligned} \tag{S24}$$

$$\begin{aligned}
U_\chi S_2^+ U_\chi^{-1} &= S_2^+ + \frac{\chi}{2} (-2S_1^+ S_{2z} I^-) + \left(\frac{\chi}{2}\right)^2 \frac{1}{2!} \left(\frac{-1}{2} S_2^+ + 2S_{1z} S_2^+ I_z\right) \\
&\quad + \left(\frac{\chi}{2}\right)^3 \frac{1}{3!} (2S_1^+ S_{2z} I^-) + \left(\frac{\chi}{2}\right)^4 \frac{1}{4!} \left(\frac{1}{2} S_2^+ - 2S_{1z} S_2^+ I_z\right) + \dots \\
&= \frac{1}{2} S_2^+ + 2S_{1z} S_2^+ I_z \\
&\quad + \left(\frac{1}{2} S_2^+ - 2S_{1z} S_2^+ I_z\right) \cos \frac{\chi}{2} - 2S_1^+ S_{2z} I^- \sin \frac{\chi}{2},
\end{aligned} \tag{S25}$$

$$\begin{aligned}
U_\chi S_2^- U_\chi^{-1} &= S_2^- + \frac{\chi}{2} (-2S_1^- S_{2z} I^+) + \left(\frac{\chi}{2}\right)^2 \frac{1}{2!} \left(\frac{-1}{2} S_2^- + 2S_{1z} S_2^- I_z\right) \\
&\quad + \left(\frac{\chi}{2}\right)^3 \frac{1}{3!} (2S_1^- S_{2z} I^+) + \left(\frac{\chi}{2}\right)^4 \frac{1}{4!} \left(\frac{1}{2} S_2^- - 2S_{1z} S_2^- I_z\right) + \dots \\
&= \frac{1}{2} S_2^- + 2S_{1z} S_2^- I_z \\
&\quad + \left(\frac{1}{2} S_2^- - 2S_{1z} S_2^- I_z\right) \cos \frac{\chi}{2} - 2S_1^- S_{2z} I^+ \sin \frac{\chi}{2}.
\end{aligned} \tag{S26}$$

Before we use the expressions above (Eq. S18 to S26), it is convenient to rewrite  $\tilde{H}_M$  (Eq. 52, S15) taking advantage of the following commutator relation:

$$\begin{aligned}
[M_{\Delta y}, S_{1z} I_z] &= [M_{\Delta y}, S_{1z} S_{2z}] = [M_{\Delta y}, S_{2z} I_z] = 0, \\
S_{1x} &= -2iS_{1y} S_{1z}, \quad S_{2x} = -2iS_{2y} S_{2z}.
\end{aligned} \tag{S27 = 54}$$

The new, more convenient form of  $\tilde{H}_M$  is therefore,

$$\begin{aligned}
\tilde{H}_M &= \frac{1}{2}(c_\alpha + c_\beta)(S_{1x} + S_{2x}) - 2i(c_\alpha - c_\beta)(S_{1y} S_{1z} I_z + S_{2y} S_{2z} I_z) \\
&\quad + 2i(s_\alpha + s_\beta)(S_{1y} - S_{2y})S_{1z} S_{2z} - 2(s_\alpha - s_\beta)(S_{1x} S_{2z} I_z - S_{2x} S_{1z} I_z).
\end{aligned} \tag{S28 = 55}$$

Now, we can evaluate the effect of the second unitary transformation  $U_\xi$ :

$$\begin{aligned}
U_\xi \tilde{H}_M U_\xi^{-1} = & (c_\alpha + c_\beta) \left\{ \frac{1}{4} (S_{1x} + S_{2x}) - (S_{1x} S_{2z} - S_{1z} S_{2x}) I_z \right. \\
& + [\frac{1}{4} (S_{1x} + S_{2x}) + (S_{1x} S_{2z} - S_{1z} S_{2x}) I_z] \cos \frac{\xi}{2} \\
& \left. + \frac{1}{2} [S_{1z} (S_2^+ I^+ + S_2^- I^-) - (S_1^+ I^- + S_1^- I^+) S_{2z}] \sin \frac{\xi}{2} \right\} \\
& + \frac{1}{2} (c_\alpha - c_\beta) \left\{ (S_{1x} + S_{2x}) I_z - (S_{1x} S_{2z} - S_{1z} S_{2x}) \right. \\
& + [(S_{1x} + S_{2x}) I_z + (S_{1x} S_{2z} - S_{1z} S_{2x})] \cos \frac{\xi}{2} \\
& \left. + \frac{1}{2} (S_2^+ I^+ + S_2^- I^- + S_1^+ I^- + S_1^- I^+) \sin \frac{\xi}{2} \right\} \\
& + \frac{1}{2} (s_\alpha + s_\beta) \left\{ (S_{1x} + S_{2x}) I_z - (S_{1x} S_{2z} - S_{1z} S_{2x}) \right. \\
& - [(S_{1x} + S_{2x}) I_z + (S_{1x} S_{2z} - S_{1z} S_{2x})] \cos \frac{\xi}{2} \\
& \left. - \frac{1}{2} (S_2^+ I^+ + S_2^- I^- + S_1^+ I^- + S_1^- I^+) \sin \frac{\xi}{2} \right\} \\
& + (s_\alpha - s_\beta) \left\{ \frac{1}{4} (S_{1x} + S_{2x}) - (S_{1x} S_{2z} - S_{1z} S_{2x}) I_z \right. \\
& - [\frac{1}{4} (S_{1x} + S_{2x}) + (S_{1x} S_{2z} - S_{1z} S_{2x}) I_z] \cos \frac{\xi}{2} \\
& \left. - \frac{1}{2} [S_{1z} (S_2^+ I^+ + S_2^- I^-) - (S_1^+ I^- + S_1^- I^+) S_{2z}] \sin \frac{\xi}{2} \right\}.
\end{aligned} \tag{S29}$$

The result of this transformation is  $\tilde{\tilde{H}}_M$ . For bookkeeping purposes, the operators in  $\tilde{\tilde{H}}_M$  can be grouped into six terms:

$$\tilde{\tilde{H}}_M = 2\omega_{1S} \cos(\omega_M t) \cdot (\tilde{\tilde{H}}_M^1 + \tilde{\tilde{H}}_M^2 + \tilde{\tilde{H}}_M^3 + \tilde{\tilde{H}}_M^4 + \tilde{\tilde{H}}_M^5 + \tilde{\tilde{H}}_M^6) \tag{S30 = 57}$$

$$\tilde{\tilde{H}}_M^1 = \tilde{\tilde{h}}_1 (S_{1x} + S_{2x}), \quad \tilde{\tilde{h}}_1 = \frac{1}{4} [(c_\alpha + c_\beta + s_\alpha - s_\beta) + (c_\alpha + c_\beta - s_\alpha + s_\beta) \cos \frac{\xi}{2}] \tag{S31}$$

$$\tilde{\tilde{H}}_M^2 = \tilde{\tilde{h}}_2 (S_{1x} + S_{2x}) I_z, \quad \tilde{\tilde{h}}_2 = \frac{1}{2} [(c_\alpha - c_\beta + s_\alpha + s_\beta) + (c_\alpha - c_\beta - s_\alpha - s_\beta) \cos \frac{\xi}{2}] \tag{S32}$$

$$\tilde{\tilde{H}}_M^3 = \tilde{\tilde{h}}_3 (S_{1x} S_{2z} - S_{1z} S_{2x}), \quad \tilde{\tilde{h}}_3 = \frac{1}{2} [(-c_\alpha + c_\beta - s_\alpha - s_\beta) + (c_\alpha - c_\beta - s_\alpha - s_\beta) \cos \frac{\xi}{2}] \tag{S33}$$

$$\tilde{\tilde{H}}_M^4 = \tilde{\tilde{h}}_4 (S_{1x} S_{2z} - S_{1z} S_{2x}) I_z, \quad \tilde{\tilde{h}}_4 = [(-c_\alpha - c_\beta - s_\alpha + s_\beta) + (c_\alpha + c_\beta - s_\alpha + s_\beta) \cos \frac{\xi}{2}] \tag{S34}$$

$$\tilde{\tilde{H}}_M^5 = \tilde{\tilde{h}}_5 (S_2^+ I^+ + S_2^- I^- + S_1^+ I^- + S_1^- I^+), \quad \tilde{\tilde{h}}_5 = \frac{1}{4} (c_\alpha - c_\beta - s_\alpha - s_\beta) \sin \frac{\xi}{2} \tag{S35}$$

$$\tilde{\tilde{H}}_M^6 = \tilde{\tilde{h}}_6 [(S_2^+ I^+ + S_2^- I^-) S_{1z} - (S_1^+ I^- + S_1^- I^+) S_{2z}], \quad \tilde{\tilde{h}}_6 = \frac{1}{2} (c_\alpha + c_\beta - s_\alpha + s_\beta) \sin \frac{\xi}{2} \tag{S36}$$



## Section V – Example 3 (Eq. 80 to 83)

Letting

$$\tilde{H}_{M3}^{eff*} = S_{1x}S_2^\alpha I^\alpha + \frac{1}{2}S_1^\alpha(S_2^+I^+ + S_2^-I^-) = S_{1x}I_{2\Sigma}^\alpha + S_1^\alpha I_{2\Sigma x}, \quad (S37 = 80)$$

we obtain

$$\begin{aligned} & [\tilde{H}_{M3}^{eff*}, S_{1z} + S_{2z}] \\ &= -iS_{1y}S_2^\alpha I^\alpha - \frac{1}{2}S_1^\alpha(S_2^+I^+ - S_2^-I^-), \end{aligned} \quad (S38 = 81)$$

$$\begin{aligned} & [\tilde{H}_{M3}^{eff*}, [\tilde{H}_{M3}^{eff*}, S_{1z} + S_{2z}]] \\ &= [\tilde{H}_{M3}^{eff*}, iS_{1y}S_2^\alpha I^\alpha + \frac{1}{2}S_1^\alpha(S_2^+I^+ - S_2^-I^-)] \\ &= [S_{1x}I_{2\Sigma}^\alpha + S_1^\alpha I_{2\Sigma x}, iS_{1y}I_{2\Sigma}^\alpha + iS_1^\alpha I_{2\Sigma y}] \\ &= i(S_{1x}I_{2\Sigma}^\alpha S_{1y}I_{2\Sigma}^\alpha + S_{1x}I_{2\Sigma}^\alpha S_1^\alpha I_{2\Sigma y} + S_1^\alpha I_{2\Sigma x}S_{1y}I_{2\Sigma}^\alpha + S_1^\alpha I_{2\Sigma x}S_1^\alpha I_{2\Sigma y} \\ &\quad - S_{1y}I_{2\Sigma}^\alpha S_{1x}I_{2\Sigma}^\alpha - S_{1y}I_{2\Sigma}^\alpha S_1^\alpha I_{2\Sigma x} - S_1^\alpha I_{2\Sigma y}S_{1x}I_{2\Sigma}^\alpha - S_1^\alpha I_{2\Sigma y}S_1^\alpha I_{2\Sigma x}) \quad (S39 = 82) \\ &= i(\frac{i}{2}S_{1z}I_{2\Sigma}^\alpha + \frac{1}{2}S_1^- \frac{1}{2i}I_{2\Sigma}^+ + \frac{1}{2i}S_1^+ \frac{1}{2}I_{2\Sigma}^- + S_1^\alpha \frac{i}{2}I_{2\Sigma z} \\ &\quad - \frac{-i}{2}S_{1z}I_{2\Sigma}^\alpha - \frac{-1}{2i}S_1^- \frac{1}{2}I_{2\Sigma}^+ - \frac{1}{2}S_1^+ \frac{-1}{2i}I_{2\Sigma}^- - S_1^\alpha \frac{-i}{2}I_{2\Sigma z}) \\ &= -S_{1z}I_{2\Sigma}^\alpha + \frac{1}{2}S_1^-I_{2\Sigma}^+ + \frac{1}{2}S_1^+I_{2\Sigma}^- - S_1^\alpha I_{2\Sigma z} \\ &= -S_1^\alpha S_2^\alpha I^\alpha + \frac{1}{2}S_1^\beta S_2^\alpha I^\alpha + \frac{1}{2}S_1^\alpha S_2^\beta I^\beta + \frac{1}{2}(S_1^+S_2^-I^- + S_1^-S_2^+I^+), \end{aligned}$$

$$\begin{aligned}
& [\tilde{\tilde{H}}_{M3}^{eff*}, [\tilde{\tilde{H}}_{M3}^{eff*}, [\tilde{\tilde{H}}_{M3}^{eff*}, S_{1z} + S_{2z}]]] \\
& = [\tilde{\tilde{H}}_{M3}^{eff*}, -S_1^\alpha S_2^\alpha I^\alpha + \frac{1}{2} S_1^\beta S_2^\alpha I^\alpha + \frac{1}{2} S_1^\alpha S_2^\beta I^\beta + \frac{1}{2} (S_1^+ S_2^- I^- + S_1^- S_2^+ I^+)] \\
& = [S_{1x} I_{2\Sigma}^\alpha + S_1^\alpha I_{2\Sigma x}, -S_{1z} I_{2\Sigma}^\alpha + \frac{1}{2} S_1^- I_{2\Sigma}^+ + \frac{1}{2} S_1^+ I_{2\Sigma}^- - S_1^\alpha I_{2\Sigma z}] \\
& = -S_{1x} I_{2\Sigma}^\alpha S_{1z} I_{2\Sigma}^\alpha + S_{1x} I_{2\Sigma}^\alpha \frac{1}{2} S_1^- I_{2\Sigma}^+ + S_{1x} I_{2\Sigma}^\alpha \frac{1}{2} S_1^+ I_{2\Sigma}^- - S_{1x} I_{2\Sigma}^\alpha S_1^\alpha I_{2\Sigma z} \\
& \quad - S_1^\alpha I_{2\Sigma x} S_{1z} I_{2\Sigma}^\alpha + S_1^\alpha I_{2\Sigma x} \frac{1}{2} S_1^- I_{2\Sigma}^+ + S_1^\alpha I_{2\Sigma x} \frac{1}{2} S_1^+ I_{2\Sigma}^- - S_1^\alpha I_{2\Sigma x} S_1^\alpha I_{2\Sigma z} \\
& \quad + S_{1z} I_{2\Sigma}^\alpha S_{1x} I_{2\Sigma}^\alpha - \frac{1}{2} S_1^- I_{2\Sigma}^+ S_{1x} I_{2\Sigma}^\alpha - \frac{1}{2} S_1^+ I_{2\Sigma}^- S_{1x} I_{2\Sigma}^\alpha + S_1^\alpha I_{2\Sigma z} S_{1x} I_{2\Sigma}^\alpha \\
& \quad + S_{1z} I_{2\Sigma}^\alpha S_1^\alpha I_{2\Sigma x} - \frac{1}{2} S_1^- I_{2\Sigma}^+ S_1^\alpha I_{2\Sigma x} - \frac{1}{2} S_1^+ I_{2\Sigma}^- S_1^\alpha I_{2\Sigma x} + S_1^\alpha I_{2\Sigma z} S_1^\alpha I_{2\Sigma x} \\
& = \frac{i}{2} S_{1y} I_{2\Sigma}^\alpha + \frac{1}{4} S_1^\alpha I_{2\Sigma}^+ + 0 - \frac{1}{4} S_1^- I_{2\Sigma}^\alpha \\
& \quad - \frac{1}{4} S_1^\alpha I_{2\Sigma}^- + 0 + \frac{1}{4} S_1^+ I_{2\Sigma}^\alpha + \frac{i}{2} S_1^\alpha I_{2\Sigma y} \\
& \quad + \frac{i}{2} S_{1y} I_{2\Sigma}^\alpha + 0 - \frac{1}{4} S_1^\alpha I_{2\Sigma}^- + \frac{1}{4} S_1^+ I_{2\Sigma}^\alpha \\
& \quad + \frac{1}{4} S_1^\alpha I_{2\Sigma}^+ - \frac{1}{4} S_1^- I_{2\Sigma}^\alpha + 0 + \frac{i}{2} S_1^\alpha I_{2\Sigma y} \\
& = 2iS_{1y} I_{2\Sigma}^\alpha + 2iS_1^\alpha I_{2\Sigma y} \\
& = 2[iS_{1y} S_2^\alpha I^\alpha + \frac{1}{2} S_1^\alpha (S_2^+ I^+ - S_2^- I^-)]. \tag{S40 = 83}
\end{aligned}$$

## Section VI – Example 4 (Eq. 89 to 92)

Letting

$$\tilde{\tilde{H}}_{M4}^{eff*} = S_{1x} S_2^\beta I^\beta + \frac{1}{2} S_1^\alpha (S_2^+ I^+ + S_2^- I^-) = S_{1x} I_{2\Sigma}^\beta + S_1^\alpha I_{2\Sigma x}, \tag{S41 = 89}$$

we obtain

$$\begin{aligned}
& [\tilde{\tilde{H}}_{M4}^{eff*}, S_{1z} + S_{2z}] \\
& = -iS_{1y} S_2^\beta I^\beta - \frac{1}{2} S_1^\alpha (S_2^+ I^+ - S_2^- I^-), \tag{S42 = 90}
\end{aligned}$$

$$\begin{aligned}
& [\tilde{\tilde{H}}_{M4}^{eff*}, [\tilde{\tilde{H}}_{M4}^{eff*}, S_{1z} + S_{2z}]] \\
& = [\tilde{\tilde{H}}_{M4}^{eff*}, iS_{1y}S_2^\beta I^\beta + \frac{1}{2}S_1^\alpha(S_2^+I^+ - S_2^-I^-)] \\
& = [S_{1x}I_{2\Sigma}^\beta + S_1^\alpha I_{2\Sigma x}, iS_{1y}I_{2\Sigma}^\beta + iS_1^\alpha I_{2\Sigma y}] \\
& = i(S_{1x}I_{2\Sigma}^\beta S_{1y}I_{2\Sigma}^\beta + S_{1x}I_{2\Sigma}^\beta S_1^\alpha I_{2\Sigma y} + S_1^\alpha I_{2\Sigma x}S_{1y}I_{2\Sigma}^\beta + S_1^\alpha I_{2\Sigma x}S_1^\alpha I_{2\Sigma y} \\
& \quad - S_{1y}I_{2\Sigma}^\beta S_{1x}I_{2\Sigma}^\beta - S_1^\alpha I_{2\Sigma y}S_{1x}I_{2\Sigma}^\beta - S_{1y}I_{2\Sigma}^\beta S_1^\alpha I_{2\Sigma x} - S_1^\alpha I_{2\Sigma y}S_1^\alpha I_{2\Sigma x}) \quad (\text{S44=} \\
& = i(\frac{i}{2}S_{1z}I_{2\Sigma}^\beta - \frac{1}{4i}S_1^-I_{2\Sigma}^- + \frac{1}{4i}S_1^+I_{2\Sigma}^+ + \frac{i}{2}S_1^\alpha I_{2\Sigma z} \\
& \quad + \frac{i}{2}S_{1z}I_{2\Sigma}^\beta - \frac{1}{4i}S_1^+I_{2\Sigma}^+ + \frac{1}{4i}S_1^-I_{2\Sigma}^- + \frac{i}{2}S_1^\alpha I_{2\Sigma z}) \\
& = -S_{1z}I_{2\Sigma}^\beta - S_1^\alpha I_{2\Sigma z} \\
& = -S_{1z}S_2^\beta I^\beta - \frac{1}{2}S_1^\alpha(S_2^\alpha I^\alpha - S_2^\beta I^\beta),
\end{aligned}$$

$$\begin{aligned}
& [\tilde{\tilde{H}}_{M4}^{eff*}, [\tilde{\tilde{H}}_{M4}^{eff*}, [\tilde{\tilde{H}}_{M4}^{eff*}, S_{1z} + S_{2z}]]] \\
& = [\tilde{\tilde{H}}_{M4}^{eff*}, S_{1z}S_2^\beta I^\beta + \frac{1}{2}S_1^\alpha(S_2^\alpha I^\alpha - S_2^\beta I^\beta)] \\
& = [S_{1x}I_{2\Sigma}^\beta + S_1^\alpha I_{2\Sigma x}, S_{1z}I_{2\Sigma}^\beta + S_1^\alpha I_{2\Sigma z}] \\
& = S_{1x}I_{2\Sigma}^\beta S_{1z}I_{2\Sigma}^\beta + S_{1x}I_{2\Sigma}^\beta S_1^\alpha I_{2\Sigma z} + S_1^\alpha I_{2\Sigma x}S_{1z}I_{2\Sigma}^\beta + S_1^\alpha I_{2\Sigma x}S_1^\alpha I_{2\Sigma z} \\
& \quad - S_{1z}I_{2\Sigma}^\beta S_{1x}I_{2\Sigma}^\beta - S_1^\alpha I_{2\Sigma z}S_{1x}I_{2\Sigma}^\beta - S_{1z}I_{2\Sigma}^\beta S_1^\alpha I_{2\Sigma x} - S_1^\alpha I_{2\Sigma z}S_1^\alpha I_{2\Sigma x} \quad (\text{S45=} \\
& = \frac{-i}{2}S_{1y}I_{2\Sigma}^\beta - \frac{1}{4}S_1^-I_{2\Sigma}^- + \frac{1}{4}S_1^+I_{2\Sigma}^+ - \frac{i}{2}S_1^\alpha I_{2\Sigma y} \\
& \quad - \frac{i}{2}S_{1y}I_{2\Sigma}^\beta + \frac{1}{4}S_1^+I_{2\Sigma}^+ - \frac{1}{4}S_1^\alpha I_{2\Sigma}^- - \frac{i}{2}S_1^\alpha I_{2\Sigma y} \\
& = -\frac{i}{2}S_{1y}I_{2\Sigma}^\beta - \frac{i}{2}S_1^\alpha I_{2\Sigma y} \\
& = -\frac{i}{2}S_{1y}S_2^\beta I^\beta - \frac{1}{4}S_1^\alpha(S_2^+I^+ - S_2^-I^-).
\end{aligned}$$

## Section VII – Experimental methods

DNP experiments were performed at 211 MHz  $^1\text{H}$  Larmor frequency (140 GHz EPR frequency) in a custom designed spectrometer (courtesy of D. J. Ruben, Francis Bitter Magnet Laboratory, Massachusetts Institute of Technology, Cambridge, MA) equipped with a home-built 4 mm low temperature MAS probe and a superconducting sweep coil. The MW source was a custom built gyrotron capable of delivering 10 W of continuous-wave MW irradiation. The samples contained 40 mM trityl (for the SE) or 10 mM

TOTAPOL (for the CE) dispersed in a glycerol-d<sub>8</sub>/D<sub>2</sub>O/H<sub>2</sub>O (60:25:15 w/w/w) matrix, also containing 2M <sup>13</sup>C-urea. Experiments were performed at 90 K, in sapphire rotors,  $\omega_r/2\pi = 5$  kHz. The enhanced <sup>1</sup>H signal of the sample was used for recording the DNP enhancement profiles and the enhancement was defined as the difference between the <sup>1</sup>H signals with MW on and off.